

Uniform-in-Time Propagation of Chaos for Consensus-Based Optimization

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Joint work with

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Consider $f : \mathbb{R}^d \rightarrow \mathbb{R}$.

Goal: Find x_* such that $f(x_*) = \min_x f(x)$.

Challenges:

- ▶ f is non-convex
- ▶ f is only given as a black-box oracle
- ▶ f is not differentiable, or derivatives are too costly
- ▶ Additional constraints may apply

Approach: Evolve a cloud of particles instead of one point.

Consider J -particles $x_1(t), \dots, x_J(t) \in \mathbb{R}^d$ satisfying

$$\dot{x}_j(t) = g_j(t, x_1(t), \dots, x_J(t)),$$

(and stochastic variants) such that $x_j(t) \rightarrow x_*$ as $t \rightarrow \infty$.

- ▶ **Goal:** Design dynamics where group self-organizes to regions of **lower function-values**
- ▶ **Features:** Parallelizable, amenable to **Mean-field Analysis** ($J \rightarrow \infty$)

Mean-field viewpoint: Analyze algorithm by tracking probability density instead of individual particles

Interacting particle methods: Gradient-based vs. Derivative-free

Gradient-based

- ▶ Stein variational gradient descent (SVGD)^a
- ▶ Swarm-Based Gradient Descent^b
- ▶ Other particle methods using ∇f or score information

Require **gradient**, **score** information to guide particle motion.

^aQ. Liu and D. Wang. In *Advances in Neural Information Processing Systems* 29, 2016.

^bJ. Lu, E. Tadmor, and A. Zenginoglu. 2022.

Derivative-free

- ▶ **Particle Swarm Optimization (PSO)**^a: classical swarm-based global optimizer
- ▶ Sequential Monte Carlo (SMC)
- ▶ Ensemble Kalman inversion (EKI)
- ▶ **Consensus-Based Optimization (CBO)**^b
- ▶ Consensus-Based Sampling (CBS)^c

Rely on **particle interaction**, **weighting**, or **function values** rather than direct gradient access.

^aJ. Kennedy and R. Eberhart. In *Proceedings of ICNN'95-international conference on neural networks*. IEEE, 1995.

^bR. Pinnau, C. Totzeck, O. Tse, and S. Martin. *Math. Models Methods Appl. Sci.*, 2017.

^cJ. A. Carrillo, F. Hoffmann, A. M. Stuart, and U. Vaes. *Stud. Appl. Math.*, 2022.

Interacting Particle System

$$dX_t^{(j)} = \underbrace{-\left(X_t^{(j)} - \mathcal{M}_\alpha(\mu_t^J)\right) dt}_{\text{drift term}} + \underbrace{\sqrt{2}\sigma \left|X_t^{(j)} - \mathcal{M}_\alpha(\mu_t^J)\right| dW_t^{(j)}}_{\text{noise term}}, \quad j = 1, \dots, J.$$

where $W_t^{(j)}$ is a d -dimensional Brownian motion, and

$$\mathcal{M}_\alpha(\mu) = \frac{\int x e^{-\alpha f(x)} \mu(dx)}{\int e^{-\alpha f(x)} \mu(dx)}, \quad \mu \in \mathcal{P}(\mathbb{R}^d).$$

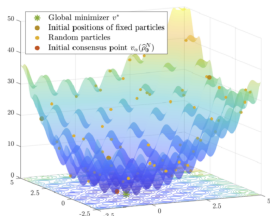
Balance of two forces:

- ▶ drift toward consensus → **exploitation**
- ▶ multiplicative noise → **exploration**

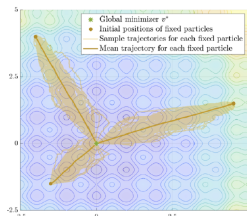
Consensus-Based Optimization (CBO)

Interacting Particle System

$$dX_t^{(j)} = \underbrace{-\left(X_t^{(j)} - \mathcal{M}_{\alpha}(\mu_t^J)\right) dt}_{\text{drift term}} + \underbrace{\sqrt{2}\sigma \left|X_t^{(j)} - \mathcal{M}_{\alpha}(\mu_t^J)\right| dW_t^{(j)}}_{\text{noise term}}, \quad j = 1, \dots, J.$$



(a) The Rastrigin function $\mathcal{E}$ and an exemplary initialization for one run of the experiment



(b) Individual agents follow, on average, the gradient flow of the map $v \mapsto \|v - v^*\|_2^2$

Figure: Excerpted from^a

^aM. Fornasier, T. Klock, and K. Riedl. *SIAM J. Optim.*, 2024.

$$dX_t^{(j)} = -\left(X_t^{(j)} - \mathcal{M}_{\alpha}(\mu_t^J)\right) dt + \sqrt{2}\sigma \left|X_t^{(j)} - \mathcal{M}_{\alpha}(\mu_t^J)\right| dW_t^{(j)}. \quad j = 1, \dots, J.$$

Advantages:

- ▶ Derivative-free
- ▶ Theoretical guarantees via mean-field analysis

Key tool: **Laplace's principle** for any $\mu \in \mathcal{P}(\mathbb{R}^d)$,

$$\mathcal{M}_{\alpha}(\mu) \rightarrow x_* \quad \text{as} \quad \alpha \rightarrow \infty \quad \text{if} \quad x_* \in \text{supp } \mu.$$

Mean-field equations:

- **SDE:** $\bar{X}_t \in \mathbb{R}^d$ solves

$$d\bar{X}_t = -\left(\bar{X}_t - \mathcal{M}_\alpha(\mu_t)\right) dt + \sqrt{2}\sigma \left|\bar{X}_t - \mathcal{M}_\alpha(\mu_t)\right| dW_t, \quad \mu_t = \text{Law}(\bar{X}_t)$$

- **PDE:** $\mu_t \in \mathcal{P}(\mathbb{R}^d)$ solves

$$\partial_t \mu_t = \nabla \cdot \left((x - \mathcal{M}_\alpha(\mu_t)) \mu_t \right) + \sigma^2 \Delta \left(|x - \mathcal{M}_\alpha(\mu_t)|^2 \mu_t \right).$$

Convergence of the mean-field solution¹²

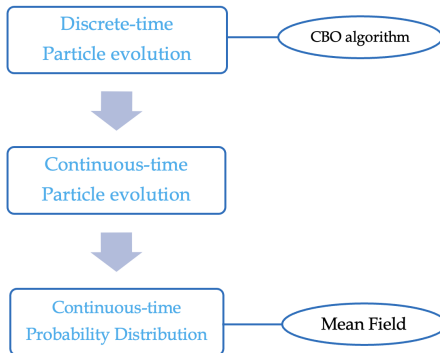
- (1) **Consensus formation:** $\mathcal{M}_\alpha(\mu_t) \rightarrow \hat{x}(\alpha)$ as $t \rightarrow \infty$ (**exponential rate**).
- (2) **Laplace's method:** $\hat{x}(\alpha) \rightarrow \arg \min_{x \in \mathbb{R}^d} f(x)$ as $\alpha \rightarrow \infty$.

¹J. A. Carrillo, Y.-P. Choi, C. Totzeck, and O. Tse. [Mathematical Models and Methods in Applied Sciences](#), 2018.

²H. Huang, J. Qiu, and K. Riedl. [Appl. Math. Optim.](#), 2023.

Limitations of Mean Field Analysis

- ▶ **Analysis** assumes a **Continuous Mean-Field PDE**
- ▶ **Algorithm** uses **Discrete-time, Finite Particles**
 - ▶ **time-discretization gap**
 - ▶ **finite-particle gap**



Paper	Result Type	Limitations
¹	Qualitative MFL	No Rate
^{2, 3}	Finite-time MFL	Time-dependent rate
⁴	Uniform-time MFL	Modified dynamics
⁵	Uniform-time Weak PoC	Modified dynamics
This Work⁶	Uniform-time Strong PoC	Original dynamics

- **Finite-in-time limits:** For $J \in \mathbb{N}$ particles,

$$\mathbb{E} \left[\sup_{0 \leq t \leq T} \left| X_t^{(j)} - \bar{X}_t \right|^p \right] \leq C J^{-\frac{p}{2}}, \text{ where } C = C(T) \rightarrow \infty \text{ as } T \rightarrow \infty$$

- **Our Result:** Remove time-dependence of C for the unmodified algorithm.

¹H. Huang and J. Qiu. *Math. Methods Appl. Sci.*, 2022.

²H. Huang, J. Qiu, and K. Riedl. *Appl. Math. Optim.*, 2023.

³N. J. Gerber, F. Hoffmann, and U. Vaes. 2024.

⁴H. Huang and H. Kouhkuh. *arXiv preprint*, 2024.

⁵E. Bayraktar, I. Ekren, and H. Zhou. *arXiv preprint*, 2025.

⁶N. Gerber, F. Hoffmann, D. Kim, and U. Vaes. 2026.

Uniform-in-time propagation of chaos [Gerber, Hoffmann, **K.**, Vaes (2025)]¹

Assume:

- ▶ f is bounded and globally Lipschitz
- ▶ $\bar{\rho}_0$ has finite moments of all orders
- ▶ Noise coefficient σ is sufficiently small

Then there exists $C_{\text{MFL}}(\alpha, \underline{f}, \bar{f}, L_f, \sigma, d)$ such that for $J \in \mathbb{N}$ particles:

$$\sup_{t \geq 0} \mathbf{E} \left[\frac{1}{J} \sum_{j=1}^J |X_t^{(j)} - \bar{X}_t^{(j)}|^2 \right] \leq \frac{C_{\text{MFL}}}{J}.$$

- ▶ Optimal Monte Carlo rate $\mathcal{O}(J^{-1})$.
- ▶ Can track how C_{MFL} depends on other parameters explicitly
- ▶ **Small-noise condition** $\Rightarrow$ **contraction towards consensus** $>$ **noise dispersion**

¹N. Gerber, F. Hoffmann, D. Kim, and U. Vaes. 2026.

Propagation of Chaos (PoC) via Coupling

Set Up

- ▶ $(W_t^j)_{j=1}^J$ are i.i.d. Brownian motions in $\mathbb{R}^d$.
- ▶ $b : \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}^d$, $\sigma : \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}$.
- ▶ $(X_t^j)_{j=1}^J$: interacting particle system, $(\bar{X}_t^j)_{j=1}^J$: i.i.d. copies of the mean-field system.

For $j = 1, \dots, J$, consider

$$X_t^j = X_0^j + \int_0^t \frac{1}{J} \sum_{k=1}^J b(X_s^j, X_s^k) ds + \int_0^t \frac{1}{J} \sum_{k=1}^J \sigma(X_s^j, X_s^k) dW_s^j,$$

(microscopic system)

$$\bar{X}_t^j = X_0^j + \int_0^t \int_{\mathbb{R}^d} b(\bar{X}_s^j, z) \rho_s(dz) ds + \int_0^t \left(\int_{\mathbb{R}^d} \sigma(\bar{X}_s^j, z) \rho_s(dz) \right) dW_s^j,$$

(i.i.d. mean-field copies)

where $\rho_t = \text{Law}(\bar{X}_t^j)$ for all $t \geq 0$.

Propagation of Chaos:

$$\mu_t^J := \frac{1}{J} \sum_{j=1}^J \delta_{X_t^j} \longrightarrow \rho_t \quad \text{as } J \rightarrow \infty.$$

Theorem (PoC via Synchronous Coupling (McKean, Sznitman^{1,2}))

If b, σ are globally Lipschitz and bounded, and $X_0^j, \bar{X}_0^j$ are sampled i.i.d. from $\bar{\rho}_0$, then:

$$\mathbb{E} [W_2^2(\mu_t^J, \rho_t)] \leq \mathbb{E} [|X_t^j - \bar{X}_t^j|^2] \leq \frac{C_1}{J} e^{\varepsilon t}$$

Synchronous Coupling Principles

- ▶ Pair particle system and J mean-field copies under **shared noise**.
- ▶ Original argument requires b, σ to be globally Lipschitz (can be adapted).
- ▶ **Limitation:** Classic estimates blow up as $t \rightarrow \infty$ ($e^{\varepsilon t}$).
- ▶ Other couplings (e.g., reflection coupling³, Optimal Coupling⁴) may lead to uniform-in-time bounds.

¹H. P. McKean. In A. K. Aziz, editor, *Lecture Series in Differential Equations, Volume 2*. Van Nostrand Reinhold Company, 1969.

²A.-S. Sznitman. *Zeitschrift für Wahrscheinlichkeitstheorie und Verwandte Gebiete*, 1984.

³A. Eberle. *Probability Theory and Related Fields*, Oct. 2015.

⁴S. Salem. 2020.

Application of Synchronous Coupling to a Toy example

Toy example (with $\mathcal{M}(\mu)$ the usual mean under μ)

Interacting particle system:

$$dX_t^j = -\left(X_t^j - \mathcal{M}(\mu_t^J)\right) dt + e^{-t} dW_t^j, \quad X_0^j = x_0^j \stackrel{\text{i.i.d.}}{\sim} \bar{\rho}_0 \quad j = 1, \dots, J.$$

Mean field limit:

$$\begin{cases} d\bar{X}_t = -\left(\bar{X}_t - \mathcal{M}(\bar{\rho}_t)\right) dt + e^{-t} d\bar{W}_t, \\ \bar{\rho}_t = \text{Law}(\bar{X}_t). \end{cases}$$

Synchronous Coupling

We couple the particle system and J copies of the mean field dynamics:

$$\begin{aligned} dX_t^j &= -\left(X_t^j - \mathcal{M}(\mu_t^J)\right) dt + e^{-t} dW_t^j, & X_0^j &= x_0^j, & j &= 1, \dots, J, \\ d\bar{X}_t^j &= -\left(\bar{X}_t^j - \mathcal{M}(\bar{\rho}_t)\right) dt + e^{-t} dW_t^j, & \bar{X}_0^j &= x_0^j, & j &= 1, \dots, J, \end{aligned}$$

with same initial condition and driving Brownian motions.

Application of Synchronous Coupling to a Toy example

$$\begin{aligned}dX_t^j &= -\left(X_t^j - \mathcal{M}(\mu_t^J)\right) dt + e^{-t} dW_t^j, & X_0^j &= x_0^j, \\d\bar{X}_t^j &= -\left(\bar{X}_t^j - \mathcal{M}(\bar{\rho}_t)\right) dt + e^{-t} dW_t^j, & \bar{X}_0^j &= x_0^j.\end{aligned}$$

Key Lemma

Global Lipschitz continuity of $\mathcal{M}$: $\mathcal{P}_1(\mathbf{R}^d) \rightarrow \mathbf{R}^d$:

$$\forall(\mu, \nu), \quad \left| \mathcal{M}(\mu) - \mathcal{M}(\nu) \right| \leq \mathcal{W}_2(\mu, \nu).$$

$$\begin{aligned}\mathbf{E} \left[\left| X_t^1 - \bar{X}_t^1 \right|^2 \right] &\lesssim \int_0^t \mathbf{E} \left| X_s^1 - \bar{X}_s^1 \right|^2 + \mathbf{E} \left| \mathcal{M}(\mu_s^J) - \mathcal{M}(\bar{\rho}_s) \right|^2 ds \\&\lesssim \int_0^t \mathbf{E} \left| X_s^1 - \bar{X}_s^1 \right|^2 + \mathbf{E} \left| \mathcal{M}(\mu_s^J) - \mathcal{M}(\bar{\mu}_s^J) \right|^2 + \mathbf{E} \left| \mathcal{M}(\bar{\mu}_s^J) - \mathcal{M}(\bar{\rho}_s) \right|^2 ds \\&\lesssim \int_0^t \mathbf{E} \left| X_s^1 - \bar{X}_s^1 \right|^2 + \mathbf{E} \left[\mathcal{W}_2(\mu_s^J, \bar{\mu}_s^J)^2 \right] ds + C_{\text{MC}} J^{-1} \\&\lesssim \int_0^t \mathbf{E} \left| X_s^1 - \bar{X}_s^1 \right|^2 ds + C_{\text{MC}} J^{-1} \quad \stackrel{\text{Grönwall}}{\rightsquigarrow} \quad \mathbf{E} \left[\left| X_t^1 - \bar{X}_t^1 \right|^2 \right] \leq C(t) J^{-1}.\end{aligned}$$

Why the classical approach fails for CBO

Synchronous coupling for CBO, $j \in \{1, \dots, J\}$

$$\begin{aligned}dX_t^j &= -\left(X_t^j - \mathcal{M}_\alpha(\mu_t^J)\right) dt + \sqrt{2}\sigma \left|X_t^j - \mathcal{M}_\alpha(\mu_t^J)\right| dW_t^j \\d\bar{X}_t^j &= -\left(\bar{X}_t^j - \mathcal{M}_\alpha(\bar{\rho}_t)\right) dt + \sqrt{2}\sigma \left|\bar{X}_t^j - \mathcal{M}_\alpha(\bar{\rho}_t)\right| dW_t^j\end{aligned}$$

Technical difficulties:

- ▶ **Non-Lipschitz drift:** The weighted mean $\mathcal{M}_\alpha: \mathcal{P}_1(\mathbf{R}^d) \rightarrow \mathbf{R}^d$ is **not globally Lipschitz**.
- ▶ **Long-time explosion:** A standard Grönwall argument yields an $\mathcal{O}(e^{Ct})$ error bound.
- ▶ **Multiplicative noise:** The diffusion term depends on the state and the empirical measure μ_t^J .

Extending Synchronous Coupling approach to CBO

Synchronous coupling for CBO, $j \in \{1, \dots, J\}$

$$\begin{aligned}dX_t^j &= -\left(X_t^j - \mathcal{M}_\alpha(\mu_t^J)\right) dt + \sqrt{2}\sigma \left|X_t^j - \mathcal{M}_\alpha(\mu_t^J)\right| dW_t^j \\d\bar{X}_t^j &= -\left(\bar{X}_t^j - \mathcal{M}_\alpha(\bar{\rho}_t)\right) dt + \sqrt{2}\sigma \left|\bar{X}_t^j - \mathcal{M}_\alpha(\bar{\rho}_t)\right| dW_t^j\end{aligned}$$

Ideas

- Requires **new tools** to avoid time-dependent blow-up.
- **Our Approach:** Classical coupling + **moment control** & **concentration inequality**

Assumptions

For simplicity, from now on we assume

- No noise ($\sigma = 0$)
- f bounded and globally Lipschitz
- Initialization in a compact set ($x_0^j \sim \bar{\rho}_0$) with $\bar{\rho}_0$ compactly supported

Extending Synchronous Coupling approach to CBO

$$\begin{aligned} \frac{d}{dt} \frac{1}{2J} \sum_{j=1}^J \mathbf{E} |X_t^j - \bar{X}_t^j|^2 &= \underbrace{-\frac{1}{J} \mathbf{E} \sum_{j=1}^J \langle X_t^j - \bar{X}_t^j, X_t^j - \mathcal{M}(\mu_t^J) - \bar{X}_t^j + \mathcal{M}(\bar{\mu}_t^J) \rangle}_{\text{(I)}} \\ &\quad - \underbrace{\frac{1}{J} \mathbf{E} \sum_{j=1}^J \langle X_t^j - \bar{X}_t^j, \mathcal{M}(\mu_t^J) - \mathcal{M}_{\alpha}(\mu_t^J) - \mathcal{M}(\bar{\mu}_t^J) + \mathcal{M}_{\alpha}(\bar{\mu}_t^J) \rangle}_{\text{(II)}} \\ &\quad - \underbrace{\frac{1}{J} \mathbf{E} \sum_{j=1}^J \langle X_t^j - \bar{X}_t^j, \mathcal{M}_{\alpha}(\bar{\rho}_t) - \mathcal{M}_{\alpha}(\bar{\mu}_t^J) \rangle}_{\text{(III)}}. \end{aligned}$$

- ▶ **(I)** good contractive term
- ▶ **(II)** weighted-mean perturbation
- ▶ **(III)** empirical sampling error

$$\frac{d}{dt} \frac{1}{2J} \sum_{j=1}^J \mathbf{E} |X_t^{(j)} - \bar{X}_t^{(j)}|^2 = \textcolor{blue}{(I)} + \textcolor{red}{(II)} + \textcolor{green}{(III)}.$$

- ▶ **(I) good contractive term:** favorable sign after pivoting around $\mathcal{M}$
- ▶ **(II) weighted-mean perturbation:** handled by a **novel stability estimate**
- ▶ **(III) empirical sampling error:** handled by **Monte Carlo estimate** + **decay of centered moments**

- ▶ Monte Carlo estimate: For any $p \geq 2$,

$$\mathbf{E} \left| \mathcal{M}_{\alpha}(\mu_{\mathcal{Z}^J}) - \mathcal{M}_{\alpha}(\pi) \right|_p^p \leq C_{\text{WM},p} \mathbf{E} \left| Z^1 - \mathbf{E} Z^1 \right|_p^p J^{-\frac{p}{2}},$$

$$\mu_{\mathcal{Z}^J} := \frac{1}{J} \sum_{j=1}^J \delta_{Z^j}, \quad \{Z^j\}_{j \in \mathbb{N}} \stackrel{\text{i.i.d.}}{\sim} \pi.$$

- ▶ Recall (decay of centered moments):

$$\mathbf{E} \left| \bar{X}_t - \mathbf{E} \bar{X}_t \right|^p \leq \mathbf{E} \left| \bar{X}_0 - \mathbf{E} \bar{X}_0 \right|^p e^{-\lambda_p t}.$$

The proof is about exploiting (I) and making (II)–(III) small uniformly in time.

Key Ingredients (1/2)

For a probability measure $\mu \in \mathcal{P}_1(\mathbf{R}^d)$, let:

$$\mathfrak{M}_p(\mu) := \int_{\mathbf{R}^d} |x - \mathcal{M}(\mu)|^p \mu(\mathrm{d}x).$$

1. **Novel Stability Estimate** for the weighted mean:

$$|\mathcal{M}_\alpha(\mu) - \mathcal{M}(\mu) - \mathcal{M}_\alpha(\nu) + \mathcal{M}(\nu)| \leq C_{\mathcal{M}} \left(\sqrt{\mathfrak{M}_2(\mu)} + \sqrt{\mathfrak{M}_2(\nu)} \right) \mathcal{W}_2(\mu, \nu).$$

2. **Decay of Centered Moments** ($\lambda_p := \lambda_p(\sigma, \alpha, \underline{f}, \bar{f})$):

$$\mathbf{E} [\mathfrak{M}_p(\mu_t^J)] \leq \mathbf{E} [\mathfrak{M}_p(\mu_0^J)] e^{-\lambda_p t}, \quad \mathbf{E} |\bar{X}_t - \mathbf{E} \bar{X}_t|^p \leq \mathbf{E} |\bar{X}_0 - \mathbf{E} \bar{X}_0|^p e^{-\lambda_p t}$$

3. Uniform-in-time **Raw Moment Control**:

$$\mathbf{E} \left[\sup_{t \geq 0} |X_t^j|^p \right]^{\frac{1}{p}} \leq C_{\text{Raw}, p} \mathbf{E} \left[|X_0^j|^p \right]^{\frac{1}{p}}, \quad \mathbf{E} \left[\sup_{t \geq 0} |\bar{X}_t|^p \right]^{\frac{1}{p}} \leq C_{\text{Raw}, p} \mathbf{E} \left[|\bar{X}_0|^p \right]^{\frac{1}{p}}$$

4. **Concentration of Measure** (Bounding Rare Excursions):

We control the probability of falling into a bad set where centered moments become too large:

$$\mathbf{P} \left[\sup_{t \geq 0} e^{\kappa t} \mathfrak{M}_2(\mu_t^J) \geq \mathfrak{M}_2(\bar{\rho}_0) + 1 \right] \lesssim J^{-\frac{q}{2}},$$
$$\mathbf{P} \left[\sup_{t \geq 0} e^{\kappa t} \mathfrak{M}_2(\bar{\mu}_t^J) \geq \mathfrak{M}_2(\bar{\rho}_0) + 1 \right] \lesssim J^{-\frac{q}{2}}.$$

Why concentration is needed

The difficult term in the stability estimate has the form

$$Q_t := \mathbf{E} \left[\left(\mathfrak{M}_2(\mu_t^J) + \mathfrak{M}_2(\bar{\mu}_t^J) \right) W_2^2(\mu_t^J, \bar{\mu}_t^J) \right].$$

- ▶ This is a **product** of moment factors and Wasserstein error
- ▶ The expectation is taken **outside** the whole product
- ▶ Average moment control alone is not enough

4. Concentration of Measure (Bounding Rare Excursions):

We control the probability of falling into a bad set where centered moments become too large:

$$\mathbf{P} \left[\sup_{t \geq 0} e^{\kappa t} \mathfrak{M}_2(\mu_t^J) \geq \mathfrak{M}_2(\bar{\rho}_0) + 1 \right] \lesssim J^{-\frac{q}{2}},$$
$$\mathbf{P} \left[\sup_{t \geq 0} e^{\kappa t} \mathfrak{M}_2(\bar{\mu}_t^J) \geq \mathfrak{M}_2(\bar{\rho}_0) + 1 \right] \lesssim J^{-\frac{q}{2}}.$$

Good set / bad set strategy

We split the expectation into:

$$Q_t = \mathbf{E}[\mathbf{1}_{\Omega_\kappa^*}(\cdots)] + \mathbf{E}[\mathbf{1}_{\Omega \setminus \Omega_\kappa^*}(\cdots)].$$

- ▶ On the Ω_κ^* (**good set**): centered moments decay + raw moments bound
- ▶ On the $\Omega \setminus \Omega_\kappa^*$ (**bad set**): concentration inequality

4. **Concentration of Measure** (Bounding Rare Excursions):

We control the probability of falling into a bad set where centered moments become too large:

$$\mathbf{P} \left[\sup_{t \geq 0} e^{\kappa t} \mathfrak{M}_2(\mu_t^J) \geq \mathfrak{M}_2(\bar{\rho}_0) + 1 \right] \lesssim J^{-\frac{q}{2}},$$
$$\mathbf{P} \left[\sup_{t \geq 0} e^{\kappa t} \mathfrak{M}_2(\bar{\mu}_t^J) \geq \mathfrak{M}_2(\bar{\rho}_0) + 1 \right] \lesssim J^{-\frac{q}{2}}.$$

5. **Monte Carlo estimate**: weighted-mean sampling error is $\mathcal{O}(J^{-p/2})$

$$\mathbf{E} \left| \mathcal{M}_{\alpha}(\mu_{\bar{\mathcal{X}}_t^J}) - \mathcal{M}_{\alpha}(\bar{\rho}_t) \right|_p^p \leq C_{\text{MC}} J^{-\frac{p}{2}}.$$

How the estimate closes

- ▶ Stability $\rightarrow$ Moment Control + Concentration inequality
- ▶ Monte Carlo estimate

Main Result 2: Stability

Uniform-in-time stability of particle dynamics

Let $(X_t^{(j)})_{j=1}^J$ and $(\tilde{X}_t^{(j)})_{j=1}^J$ solve the CBO system with the **same Brownian motions** but different initial data. Then for all $q \geq 2$, for small enough σ ,

$$\sup_{t \geq 0} \mathbf{E} \left[\frac{1}{J} \sum_{j=1}^J |X_t^{(j)} - \tilde{X}_t^{(j)}|^2 \right] \leq C_{\text{Stab},1} \mathbf{E} \left[\frac{1}{J} \sum_{j=1}^J |X_0^{(j)} - \tilde{X}_0^{(j)}|^2 \right] + \frac{C_{\text{Stab},2}}{J^q}.$$

Proof features:

- ▶ Same proof architecture as PoC
 - ▶ stability estimate, contractivity, moment bounds, concentration inequalities
- ▶ Monte Carlo term is not needed (particle-to-particle)
 - ▶ Residual error decays faster: J^{-q}

Main takeaway

- ▶ **Quantitative, Uniform-in-time** bridge between finite-particle CBO and its mean-field model
- ▶ Works for the **original, unmodified dynamics**
- ▶ **Key challenge resolved**: Uniform-in-time **Moment Control**
- ▶ Leveraged contractivity and stochastic analysis (concentration, BDG)

Next questions

- ▶ Can we close the **time-discretization gap** uniformly in time?
- ▶ Can we analyze **Kinetic CBO** (second-order model connection to Particle Swarm Optimization)? **(on arXiv soon)**
- ▶ Can we extend this to Consensus-Based Sampling (CBS)?
- ▶ Can we study the method with **adaptive α** ?
- ▶ Can we unify the framework for derivative-free optimization algorithms?

Acknowledgements



(a) Prof. Seung-Yeal Ha



(b) Prof. Franca Hoffmann



(c) Prof. Urbain Vaes



(d) Nicolai Gerber

Figure: Mentors and Collaborators

Thank you for your attention!

For questions, contact [Dohyeon Kim](#) via dohyeon@caltech.edu



E. Bayraktar, I. Ekren, and H. Zhou. Uniform-in-time weak propagation of chaos for consensus-based optimization. [arXiv preprint, 2502.00582](#), 2025.



J. A. Carrillo, Y.-P. Choi, C. Totzeck, and O. Tse. An analytical framework for consensus-based global optimization method. [Mathematical Models and Methods in Applied Sciences](#), 28(6):1037–1066, 2018.



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Laplace's method can be employed for studying the limit as $\beta \rightarrow \infty$ of the integral

$$I_\beta(\varphi) = \frac{\int_{\mathbf{R}^d} \varphi(x) e^{-\beta J(x)} \mu(dx)}{\int_{\mathbf{R}^d} e^{-\beta J(x)} \mu(dx)} =: \int_{\mathbf{R}^d} \varphi d(\mathcal{R}_\beta \mu), \quad \mathcal{R}_\beta : \mu \mapsto \frac{\mu e^{-\beta J}}{\int \mu e^{-\beta J}}.$$

Let $x_* = \arg \min J$. Under appropriate assumptions, it holds

$$I_\beta(\varphi) = \int_{\mathbf{R}^d} \varphi dg_\beta + \mathcal{O}\left(\frac{1}{\beta^2}\right) \quad \text{as } \beta \rightarrow \infty.$$

where $g_\beta = \mathcal{N}\left(x_*, \beta^{-1}(\text{Hess } J(x_*))^{-1}\right)$. In other words $\mathcal{R}_\beta \mu \approx g_\beta$ for large β .

Motivation:

$$e^{-\beta J(x)} \approx e^{-\beta \left(J(x_*) + \frac{1}{2} \text{Hess } J(x_*) : ((x-x_*) \otimes (x-x_*)) \right)}$$

From Implementation to Mean Field Limit [APPENDIX]

Particle-level, Discrete time: $x_h^{(j)} \in \mathbb{R}^d$ for $j = 1, \dots, J$ solves

$$x_h^{(j)}(n+1) = x_h^{(j)}(n) - h \left(x_h^{(j)} - \mathcal{M}_{\beta}(X_h) \right) + \sqrt{2h}\sigma \left| x_h^{(j)} - \mathcal{M}_{\beta}(X_h) \right| Z_h^{(j)}(n),$$

where

$$\mathcal{M}_{\beta}(X_h) = \frac{\sum_{k=1}^J x_h^{(k)} \exp(-\beta J(x_h^{(k)}))}{\sum_{k=1}^J \exp(-\beta J(x_h^{(k)}))}, \quad X_h := \left(x_h^{(j)} \right)_{j=1}^J \in \mathbb{R}^{dJ}.$$

Particle-level, Continuous time: $X_t^{(j)} \in \mathbb{R}^d$ for $j = 1, \dots, J$ solves

$$dX_t^{(j)} = - \left(X_t^{(j)} - \mathcal{M}_{\beta}(\mu_t^J) \right) dt + \sqrt{2}\sigma \left| X_t^{(j)} - \mathcal{M}_{\beta}(\mu_t^J) \right| dW_t^{(j)},$$

where

$$\mathcal{M}_{\beta}(\mu_t^J) = \frac{\int x e^{-\beta f(x)} \mu_t^J(dx)}{\int e^{-\beta f(x)} \mu_t^J(dx)} = \frac{\sum_{k=1}^J X_t^{(k)} \exp(-\beta f(X_t^{(k)}))}{\sum_{k=1}^J \exp(-\beta f(X_t^{(k)}))}, \quad \mu_t^J = \frac{1}{J} \sum_{k=1}^J \delta_{X_t^{(k)}}.$$

Mean-field limit: $\bar{x} \in \mathbb{R}^d$ solves

$$d\bar{x}_t = - \left(\bar{x}_t - \mathcal{M}_{\beta}(\mu_t) \right) dt + \sqrt{2}\sigma \left| \bar{x}_t - \mathcal{M}_{\beta}(\mu_t) \right| dW_t,$$

where

$$\mathcal{M}_{\beta}(\mu_t) = \frac{\int_{\mathbb{R}^d} \bar{x} e^{-\beta f(\bar{x})} \mu(d\bar{x})}{\int_{\mathbb{R}^d} e^{-\beta f(\bar{x})} \mu(d\bar{x})}, \quad \mu = \text{Law}(\bar{x}_t) \in \mathcal{P}(\mathbb{R}^d).$$

Proof Sketch

Define $\alpha_t^j := X_t^j - \bar{X}_t^j$.

$$\frac{d}{dt} \mathbb{E} \left[\frac{1}{2} |\alpha_t^j|^2 \right] \leq \text{Term (A)} + \text{Term (B)}$$

$$\text{Term (A)} := \frac{1}{J} \sum_{k=1}^J \mathbb{E}[\langle b(X_t^j, X_t^k) - b(\bar{X}_t^j, \bar{X}_t^k), \alpha_t^j \rangle]$$

$$\text{Term (B)} := \mathbb{E} \left[\left\langle \frac{1}{J} \sum_{k=1}^J b(\bar{X}_t^j, \bar{X}_t^k) - \int b(\bar{X}_t^j, z) \rho_t(dz), \alpha_t^j \right\rangle \right]$$

► Use Young's inequality:

$$\text{Term (B)} \leq \frac{1}{2\varepsilon} \mathbb{E}[|\alpha_t^j|^2] + \frac{\varepsilon}{2} \mathbb{E} \left[\left| \frac{1}{J} \sum_{k=1}^J b(\bar{X}_t^j, \bar{X}_t^k) - \int b(\bar{X}_t^j, z) \rho_t(dz) \right|^2 \right]$$

► The second part is $\mathcal{O}(1/J)$ by Law of Large Numbers.

Key ingredients

For $\mu \in \mathcal{P}_1(\mathbb{R}^d)$, define the centered moment

$$\mathfrak{M}_p(\mu) := \int_{\mathbb{R}^d} |x - \mathcal{M}(\mu)|^p \mu(\mathrm{d}x).$$

1. Stability estimate for the weighted mean

$$|\mathcal{M}_{\alpha}(\mu) - \mathcal{M}(\mu) - \mathcal{M}_{\alpha}(\nu) + \mathcal{M}(\nu)| \leq C_{\mathcal{M}} \left(\sqrt{\mathfrak{M}_2(\mu)} + \sqrt{\mathfrak{M}_2(\nu)} \right) W_2(\mu, \nu)$$

2. Decay of centered moments and uniform raw moment bounds

3. **Concentration Inequalities:** bad events where moments get too large have probability $\lesssim J^{-q/2}$

4. **Monte Carlo estimate:** weighted-mean sampling error is $\mathcal{O}(J^{-p/2})$

Takeaway

- ▶ Bypasses non-Lipschitz issues
- ▶ Large errors occur with vanishingly small probability

Theorem (Burkholder–Davis–Gundy inequality)

Let $(W_t)_{t \geq 0}$ denote a standard Brownian motion in $\mathbf{R}^m$ and let $(\mathcal{F}_t)_{t \geq 0}$ be the induced filtration. Let $(g_t)_{t \geq 0}$ be a $\mathbf{R}^{n \times m}$ -valued $\mathcal{F}_t$ -adapted process such that for every time $T \geq 0$, it holds that $\int_0^T \|g(t)\|_{\mathbf{F}}^2 dt < +\infty$ almost surely. Denote

$$X_t := \int_0^t g(s) dW_s \quad \text{and} \quad \langle X \rangle_t := \int_0^t \|g(s)\|_{\mathbf{F}}^2 ds.$$

Then for all $p > 0$, there exist positive constants $c_{\text{BDG},p}, C_{\text{BDG},p} < +\infty$ such that

$$\forall t \geq 0, \quad c_{\text{BDG},p} \mathbb{E} \left[\langle X \rangle_t^{\frac{p}{2}} \right] \leq \mathbb{E} \left[\sup_{0 \leq s \leq t} |X_s|^p \right] \leq C_{\text{BDG},p} \mathbb{E} \left[\langle X \rangle_t^{\frac{p}{2}} \right]. \quad (5.1)$$

Particle Swarm Optimization (PSO)

$$\begin{aligned} dx_i(t) &= v_i(t)dt, \\ mdv_i(t) &= -\gamma v_i(t)dt + \lambda_1(y_i(t) - x_i(t))dt + \lambda_2(\bar{y}(t) - x_i(t))dt \\ &\quad + \sigma_1 S(y_i(t) - x_i(t))dW_i^2(t) + \sigma_2 S(\bar{y}(t) - x_i(t))dW_i^2(t). \end{aligned}$$

- ▶ $y_i(t)$: local best (personal best position for each $x_i(t)$),
 $\bar{y}(t)$: global best (min among all y_i 's)
- ▶ 1) Eliminate local best (memory) effect
- ▶ 2) Regularize $\bar{y}$ via softmin ($\mathcal{M}_\beta$)
→ Second order CBO (: **Kinetic CBO**)
- ▶ 3) Zero inertia limit ($m \rightarrow 0$)
→ First order CBO
- ▶ Despite its wide implementation, its convergence analysis is yet nascent

Let $S : \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}$ is a Lipschitz function with Lipschitz constant L .

Particle Evolution

Let given J particles and their velocity vectors $(x_j(t), v_j(t))_{j=1}^J$ solve

$$\begin{aligned} dx_j(t) &= v_j(t)dt, \\ mdv_j(t) &= -\gamma v_j(t)dt - (x_j(t) - \mathcal{M}_{\beta}(\mu_t^J))dt + \sigma S(x_j(t) - \mathcal{M}_{\beta}(\mu_t^J))dW_j(t), \end{aligned}$$

Liouville Equation

$$\begin{aligned} \partial_t f_J &= - \sum_{j=1}^J \nabla_{x_j} \cdot (v_j f_J) + \frac{1}{m} \nabla_{v_j} \cdot \left((\gamma v_j + x_j - \mathcal{M}_{\beta}(f_J)) f_J \right) \\ &\quad + \frac{\sigma^2}{2m^2} \nabla_v \cdot \nabla_v \left(S(x_j - \mathcal{M}_{\beta}(f_J))^2 f_J \right). \end{aligned}$$

Mean Field Limit SDE and PDE

When $J \rightarrow \infty$,

$$\begin{aligned} d\bar{X}(t) &= \bar{V}(t), \\ m d\bar{V}(t) &= -\gamma V(t) - (\bar{X}(t) - \mathcal{M}_{\beta}(\mu_t))dt + \sigma S(\bar{X}(t) - \mathcal{M}_{\beta}(\mu_t))dW(t). \end{aligned} \quad (5.2)$$

The corresponding PDE to (5.2) is

$$\partial_t f + v \cdot \nabla_x f - \frac{1}{m} \nabla_v \cdot ((\gamma v + x - \mathcal{M}_{\beta}(f)) f) = \frac{\sigma^2}{2m^2} \nabla_v \cdot \nabla_v (S(x - \mathcal{M}_{\beta}(f))^2 f). \quad (5.3)$$